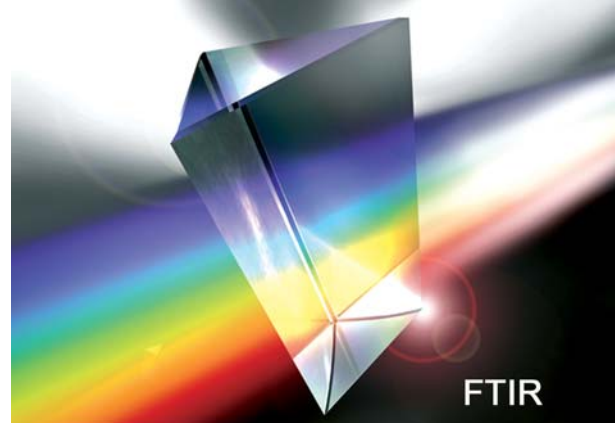


# Application Note

## Kramers-Kronig Transform and Applications



Spectral measurement can be made easily using specular spectrometry without destroying the sample. However, when it is applied to the measurement of a sample which is highly specular like a glassy or crystalline substance and which exhibits absorption in the infrared region, abnormal dispersion occurs, making the peaks of the spectrum distorted to look like those of a first derivative curve. From such a spectrum, analysis and identification of functional groups is difficult. Therefore this distorted spectrum needs to be transformed into an ordinary spectrum for interpretation. The Kramers-Kronig transform is used for that purpose. Given here is an explanation of the Kramers-Kronig transform and its applications.

### Kramers-Kronig Transform

Assuming the infrared radiation is incident normal (90°) to the surface of the sample, when the complex index of refraction ( $n^*$ ) of the substance is  $n^* = n + ik$  ( $n, k$ : refractive index and absorption coefficient of the substance), the reflectivity amplitude  $r$  for vertically incident light and the reflectivity energy  $R$  that is directly obtained from the specular reflectance spectrum are expressed as follows.

$$r = \sqrt{R} e^{i\phi} = \sqrt{R} (\cos\phi + i \sin\phi) = \frac{n - ik - 1}{n - ik + 1} \quad (1)$$

$$R = |r|^2 = r \cdot r^* = \frac{(n-1)^2 + k^2}{(n+1)^2 + k^2} \quad (2)$$

Here,  $\phi$  represents the phase change from the absorbance that occurs with reflection from the surface of a substance. The term  $r^*$  is the complex conjugate of  $r$ . When  $n$  and  $k$  in equation (1), which includes a real part and an imaginary part, are solved for, the following equations are obtained.

$$n = \frac{1 - R}{1 + R - 2\sqrt{R} \cos\phi} \quad (3)$$

$$K = \frac{-2\sqrt{R} \sin\phi}{1 + R - 2\sqrt{R} \cos\phi} \quad (4)$$

The logarithmic expression of equation (1) is as follows:

$$\ln r = \ln \sqrt{R} + i\phi \quad (5)$$

In equation (5),  $\sqrt{R}$  and  $\phi$  are mutually dependent according to the Kramers-Kronig equation:

$$\phi(v_g) = \frac{2v_g}{\pi} \int_0^\infty \frac{\ln \sqrt{R}(v)}{v^2 - v_g^2} dv \quad (6)$$

Using this equation, the phase change at an arbitrary wavenumber  $v_g$  can be calculated if the reflection energy can be measured over the entire wavenumber range, and then from equations (3) and (4), the optical constants  $n$  and  $k$  can be obtained. From such a spectrum, analysis and identification of functional groups is difficult. Therefore this distorted spectrum needs to be transformed into an ordinary spectrum for interpretation. The Kramers-Kronig transform is used for that purpose. Given here is an explanation of the Kramers-Kronig transform and its applications. and  $k$  can be obtained. Consequently, if  $k$  is calculated for the middle infrared region, e.g., 4600 - 400  $\text{cm}^{-1}$ , at a given wavenumber interval, an absorption coefficient spectrum equivalent to a transmission spectrum can be obtained from the specular reflectance spectrum. The integral (6) has  $v = v_g$  as a limit. Several integration methods have been proposed, among which Maclaurin's method and the double Fourier transformation are the most commonly applied. By Maclaurin's method, phase change ( $v_g$ ) is given by the equation (7) wherein the lower limit  $v_j$  is set so that  $v = v_g$  will not occur, and every other data point is utilized. While, by the double Fourier transformation,  $\phi(v_g)$  is obtained as an approximation of equation (6) through the double Fourier transform, as in equation (8). The Maclaurin method features higher calculation accuracy, but as it usually takes a relatively longer time for calculation, the double Fourier transform is used.

$$\phi(v_g) = \frac{2v_g}{\pi} \times 2h \times \sum_j \frac{\ln \sqrt{R}(v_j)}{v_j^2 - v_g^2} \quad (7)$$

where  $h = v_{j+1} - v_j$  and if data interval  $g$  is an odd number then  $j = 2, 4, 6, \dots, g-1, g+1, \dots$ , while if  $g$  is an even number then  $j = 1, 3, 5, \dots, g-1, g+1, \dots$

$$\phi(v_g) = 4 \int_0^\infty \cos(2 \pi v g t) dt \int_0^\infty \ln \sqrt{R(v)} \sin(2 \pi v t) dv \quad (8)$$

Table 1 Analytical Conditions of Fig. 1 and Fig. 3

Resolution : 4 cm<sup>-1</sup>

Accumulation : 100

Apodization : HappGenzel

Mirror Speed : 2.0 mm/sec

Detector : Pyroelectric Detector

## Application to Black Rubber

In the infrared spectrometry of black rubber containing carbon black, the **ATR** method with a Ge prism is applied in general. But, in the low wavenumber range below 650 cm<sup>-1</sup>, Ge itself absorbs infrared rays, which interferes with the sample absorption spectrum. In the case of a lustrous or shiny black rubber, the infrared spectrum can be measured up to 650 cm<sup>-1</sup> easily by acquiring the specular reflectance spectrum and converting it with the Kramers-Kronig transform. Fig. 1 shows the specular reflectance spectrum of a black fluorinated rubber. Under these conditions, the spectrum is distorted, which makes it difficult to obtain any qualitative information. From the spectrum after the Kramers-Kronig transform, as shown in Fig. 2, it can be confirmed that the sample is a fluorinated rubber.

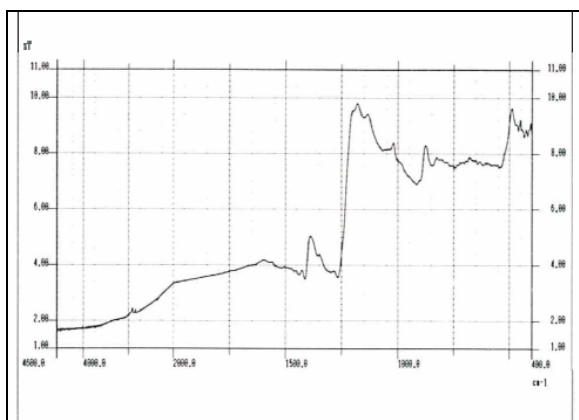


Fig. 1 Specular Reflectance Spectrum of Black Fluorinated Rubber

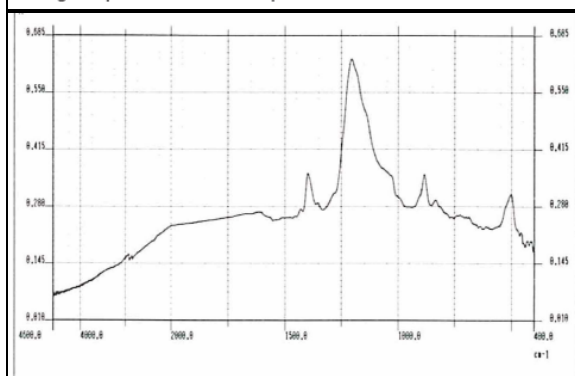


Fig. 2 Spectrum of Black Fluorinated Rubber after Kramers-Kronig Transformation

## Application to Plastic Slab

Fig. 3 shows a specular reflectance spectrum of a polymethylmethacrylate (**PMMA**) slab of a thickness of 1 cm. In the case of such a thick sample, the usual **ATR** method would require some preparation of the sample, but the specular reflectance measurement enables spectrum acquisition to be easily conducted without the need for any preparation.

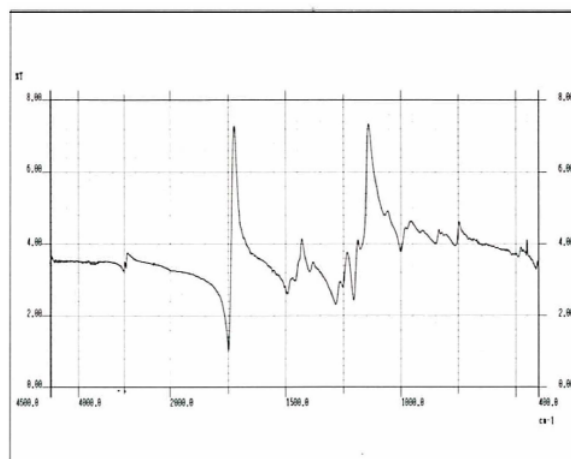


Fig. 3 Specular Reflectance Spectrum of Polymethylmethacrylate (PMMA) Slab

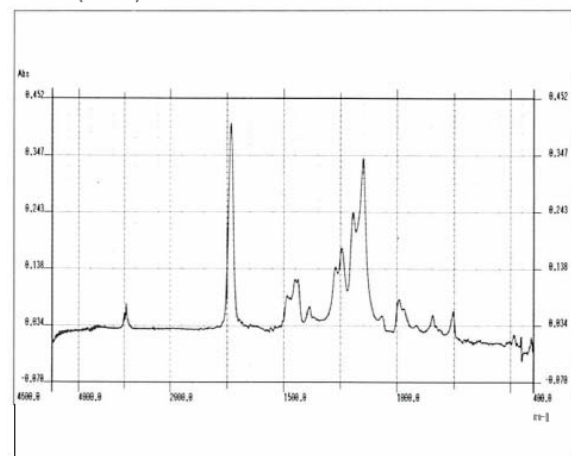


Fig. 4 Spectrum of PMMA Slab after Kramers-Kronig Transformation.

Reference: Shimadzu C103-E031 No. 228